On approximation by exponentials

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Abstract We consider an approximation of $L^2[0,\infty)$ functions by linear combinations of exponentials $\{\exp(-\lambda_\ell t)\}$. Having derived explicitly by Fourier transform techniques an orthogonal basis of exponentials, we specialize the discussion to the choice $\lambda_\ell = q^\ell$, $\ell = 0, 1, \ldots$, where $q \in (0, 1)$. In that case the underlying orthogonal functions possess a particularly appealing form and they obey interesting recurrence relations. We conclude the paper with a brief discussion of convergence issues.

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1. Orthogonality by sums of exponentials

Let λ be a complex number with positive real part and define the function $g_{\lambda} : \mathbb{R} \to \mathbb{R}$ by

$$g_{\lambda}(x) = \begin{cases} e^{-\lambda x}, & x \ge 0, \\ 0, & x < 0 \end{cases}$$
(1.1)

Thus g_{λ} is a member of the Hilbert space $H := \{f \in L^2(\mathbb{R}^2) : \operatorname{supp} f \subset [0,\infty)\}$, endowed with the usual inner product. Further, for any $f \in H$ we have the useful relation

$$(f, g_{\lambda}) = f(-i\lambda), \tag{1.2}$$

where we have used the well known fact (see Dym and McKean (1972)) that the Fourier transform of any member of H can be analytically continued throughout the lower half plane $\{z \in \mathbb{C} : \text{Im } z < 0\}$. We shall also need the Fourier transform of g_{λ} :

$$\hat{g}_{\lambda}(z) = \frac{-\mathrm{i}}{z - \mathrm{i}\lambda}.\tag{1.3}$$

Now let $(\lambda_k)_{k=0}^{\infty}$ be any sequence of complex numbers with positive real parts. Our aim is to study the subspaces

$$H_n = \operatorname{span} \{ g_{\lambda_0}, g_{\lambda_1}, \dots, g_{\lambda_n} \}$$
(1.4)

The simple form of (1.3) allows us to construct orthogonal bases for these subspaces. **Theorem 1** $\hat{H}_n \ominus \hat{H}_{n-1}$ is the one-dimensional subspace generated by the function

$$\widehat{r}_n(z) = \frac{-\mathrm{i}}{z - \mathrm{i}\lambda_n} \prod_{k=0}^{n-1} \frac{z + \mathrm{i}\lambda_k}{z - \mathrm{i}\lambda_k}.$$
(1.5)

Proof Every element of $\hat{H}_n \ominus \hat{H}_{n-1}$ is a linear combination of the rational functions $\{(z - i\lambda_k)^{-1} : 0 \le k \le n\}$; this defines \hat{r}_n up to multiplication by a constant. \Box

Partial fraction decomposition of the rational function (1.5) yields the coefficients of r_n when expressed as a linear combination of $g_{\lambda_0}, \ldots, g_{\lambda_n}$.

Proposition 2 We have

$$r_n(t) = \sum_{j=0}^n \exp(-\lambda_j t) \frac{\prod_{k=0}^{n-1} (\lambda_j + \lambda_k)}{\prod_{\ell=0, \ell \neq j}^n (\lambda_j - \lambda_\ell)}.$$
 (1.6)

Proof Writing $\hat{r}_n(z) = -i \sum_{k=0}^n \alpha_k (z - i\lambda_k)^{-1}$, we obtain

$$-\mathrm{i}\sum_{k=0}^{n}\alpha_{k}\prod_{\ell=0,\ell\neq k}^{n}(z-\mathrm{i}\lambda_{\ell})=-\mathrm{i}\prod_{k=0}^{n-1}(z+i\lambda_{k}).$$

Setting $z = i\lambda_i$ provides the relation

$$\alpha_j \prod_{\ell=0, \ell \neq j}^n (\lambda_j - \lambda_\ell) = \prod_{k=0}^{n-1} (\lambda_j + \lambda_k),$$

whence the result.

The formula derived in (1.6) is easily recognized to be a certain divided difference. **Proposition 3** Let $s_n(z,t) = \exp(-zt) \prod_{k=0}^{n-1} (z+\lambda_k)$. Then

$$r_n(t) = s_n(\cdot, t)[\lambda_0, \lambda_1, \dots, \lambda_n].$$
(1.7)

Proof The standard algebraic identity

$$f[\lambda_0, \dots, \lambda_n] = \sum_{j=0}^n \frac{f(\lambda_j)}{\prod_{\ell=0, \ell \neq j}^n (\lambda_j - \lambda_\ell)}$$

is valid for every function f defined at the points $\lambda_0, \ldots, \lambda_n$.

As a consequence of this divided difference relation, we find that the normalization chosen for r_n implies the equation $r_n(0) = 1$.

Corollary 4 For every n, we have $r_n(0) = 1$.

Proof When t = 0, the function $z \mapsto s_n(z, 0)$ is a monic polynomial of degree n. Hence the divided difference appearing on the right hand side of (1.7) is equal to one for any choice of $\lambda_0, \ldots, \lambda_n$.

Corollary 5 The two-norm of r_n is given by $||r_n||^2 = (2\lambda_n)^{-1}$.

Proof The Parseval Theorem provides the equation

$$||r_n||^2 = (2\pi)^{-1} \int_{\mathbb{R}} |\hat{r}_n(z)|^2 \, \mathrm{d}z.$$

However, for real x we have $|\hat{r}_n(x)|^2 = (x^2 + \lambda_n^2)^{-1}$. Therefore we need only compute the elementary integral

$$||r_n||^2 = (2\pi)^{-1} \int_{\mathbb{R}} (x^2 + \lambda_n^2)^{-1} \, \mathrm{d}x = (2\lambda_n)^{-1}.$$

The complex analytic theory of the Müntz theorem is closely related to the material discussed here. For example, if $(\lambda_k)_{k=0}^{\infty}$ is a sequence of positive numbers possessing a convergent subsequence with positive limit, then linear combinations of the functions $g_{\lambda_0}, g_{\lambda_1}, \ldots$ are dense in H. For, suppose $f \in H$ were orthogonal to these functions. The the Fourier transform \hat{f} satisfies $\hat{f}(-i\lambda_k) = 0$ for all k, and, by the principle of

isolated zeros for analytic functions, must therefore vanish identically. The Parseval theorem

$$\int_{\mathbb{R}} |f(x)|^2 \, \mathrm{d}x = (2\pi)^{-1} \int_{\mathbb{R}} |\hat{f}(z)|^2 \, \mathrm{d}z$$

then implies that f vanishes almost everywhere. The particular choice $\lambda = \alpha + q^k$, where 0 < q < 1 and $\alpha > 0$, obviously accumulates at the point α , and this particular case is extensively studied below.

The proof strategy of the preceding paragraph yields a derivation of Lerch's uniqueness theorem for Laplace tranforms that deserves to be better known.

Theorem 6 Let $f: [0, \infty) \to \mathbb{R}$ be a measurable function for which

$$\int_0^\infty \exp(-s_0 t) |f(t)| \, \mathrm{d}t < \infty$$

for some $s_0 \geq 0$, and the Laplace transform

$$\int_{0}^{\infty} e^{-st} f(t) dt = 0, \qquad (1.8)$$

for all sufficiently large s. Then f vanishes almost everywhere.

Proof The function $g(t) := \exp(-s_0 t)f(t)$ is absolutely integrable. Hence the dominated convergence theorem implies the continuity of its Fourier transform

$$\hat{g}(z) = \int_{0}^{\infty} e^{-izt} g(t) dt$$

for Im z < 0. Applying Morera's theorem, we deduce that the Fourier transform is, in fact, analytic for Im z < 0. However, 1.8 implies that \hat{g} vanishes on an infinite subinterval of the imaginary axis, and must therefore vanish everywhere by the principle of isolated zeros. Hence f vanishes almost everywhere.

We have not seen this proof in the literature, but its novelty is implausible.

The Müntz theorem is often proved using Cauchy's determinant identity; see, for instance, Lemma 11.3.1 of Davis (1975). It is noteworthy that our construction of r_n enables us to bypass this identity in an illuminating way. Specifically, let $g_p(t) = \exp(-pt)$, where p is not one of the numbers in the sequence $\{\lambda_k\}$. Then we can explicitly determine the distance from g_p to H_n .

Proposition 7 We have

$$\operatorname{dist}(g_p, H_n)^2 = (2p)^{-1} \left(\prod_{k=0}^n \frac{p - \lambda_k}{p + \lambda_k} \right)^2$$

Proof Replacing λ_{n+1} by p protem in Theorem 1, we see that the closest function $f_n \in H_n$ to g_p must satisfy the relation

$$\hat{g}_p(z) - \hat{f}_n(z) = \frac{-\mathrm{i}c}{z - \mathrm{i}p} \prod_{k=0}^n \frac{z + \mathrm{i}\lambda_k}{z - \mathrm{i}\lambda_k},$$

where the constant c is chosen so that the coefficient of g_p is unity. Applying (1.6), we conclude that

$$c = \prod_{k=0}^{n} \frac{p - \lambda_k}{p + \lambda_k},$$

and Corollary 5 completes the proof.

It is an interesting elementary exercise to deduce the Cauchy determinant identity from Proposition 7. To deduce one half of the Müntz theorem, let $0 < \lambda_0 < \lambda_1 < \cdots$ be any sequence for which $\lambda_k \to \infty$ and $\sum \lambda_k^{-1}$ is finite. Then the infinite product

$$\Delta(z) = \prod_{k=0}^{\infty} \frac{1 - z/\lambda_k}{1 + z/\lambda_k}$$
(1.9)

is absolutely convergent. Thus Δ is an analytic function in, say, the domain $C \setminus (-\infty, 0]$ whose only zeros are located at the points $\{\lambda_k\}$. But we have the relation

$$\lim_{n \to \infty} 2p \operatorname{dist}(g_n, H_n) = \Delta(p)^2 > 0.$$

We conclude this section by presenting a universal differential recurrence relation which is obeyed by members of the sequence $\{r_n\}_n^\infty$. Our point of departure is (1.5), which immediately implies that

$$(iz + \lambda_n)\hat{r}_n(z) = (iz - \lambda_{n-1})\hat{r}_{n-1}(z),$$
 (1.10)

whence

$$\hat{r'}_{n-1} - \hat{r'}_{n-1} = \lambda_{n-1}\hat{r}_{n-1} + \lambda_n\hat{r}_n.$$

Therefore, since the Fourier transform is an isometric linear isomorphism on $L^2(\mathbb{R})$, and by virtue of the analyticity of each r_n , we deduce

Theorem 8 The sequence r_0, r_1, \ldots obeys the differential recurrence relation

$$r'_{n-1} - r'_n = \lambda_{n-1} r_{n-1} + \lambda_n r_n, \qquad n = 1, 2, \dots$$
 (1.11)

Formula (1.11) can be recast into an interesting form. We commence by noting that (1.5) implies the recursion

$$\hat{r}_n(z) = \left(\frac{z + i\lambda_{n-1}}{z - i\lambda_n}\right) \hat{r}_{n-1}(z)$$

= $\hat{r}_{n-1}(z) - (\lambda_{n-1} + \lambda_n) \widehat{g_{\lambda_n}} r_{n-1}(z)$

Thus

$$r_n(t) = r_{n-1}(t) - (\lambda_{n-1} + \lambda_n) \int_0^t e^{-\lambda_n(t-\tau)} r_{n_1}(\tau) \,\mathrm{d}\tau.$$
(1.12)

2. Approximation by exponentials with rescaling

A particularly appealing choice of parameters is $\lambda_k = q^k + \alpha$, where $q \in (0, 1)$ and α is positive. According to Section 1, the functions $\exp(-q^k x - \alpha)$, $k = 0, 1, \ldots$, are dense with respect to the inner product

$$\langle f,g\rangle = \int_0^\infty f(x)g(x)\,\mathrm{d}x$$

or, alternatively, $\exp(-q^k x)$, $k = 0, 1, \ldots$, are dense with respect to

$$\langle f, g \rangle_{\alpha} = \int_0^\infty f(x) g(x) \mathrm{e}^{-2\,\alpha x} \,\mathrm{d}x.$$
 (2.1)

Moreover, the linear spaces H_n which have been defined in (1.4) are closed under shifts

 $f \in H_n \implies f(\cdot + \beta) \in H_n \text{ for all } \beta \in \mathbb{R};$

whilst a dilation of the independent variable by a factor of q moves up the chain $H_0 \subset H_1 \subset H_2 \subset \cdots$

$$f \in H_n \implies f(q \cdot) \in H_{n+1}.$$

We use (1.6) to describe an orthogonal basis of $+_{n=0}^{\infty} H_n$ in a closed form,

$$r_m^{[\alpha]}(t) = e^{-\alpha t} \sum_{j=0}^m \frac{\prod_{k=0}^{m-1} (q^k + q^j + 2\alpha)}{\prod_{\ell=0, \ell \neq j}^m (q^j - q^\ell)} e^{-q^j t}, \qquad m = 0, 1, \dots.$$
(2.2)

The last expression can be somewhat simplified by using Gauß-Heine symbols but this procedure has not led to significant additional insight. However, a substantially simplified form, accompanied by a wealth of further results – generating functions, recurrence relations, connections to certain functional-differential equations – follows in the case $\alpha = 0$. The *quid pro quo* is, of course, that density with respect to the inner product (2.2) is lost. Although it is not difficult to prove that density is retained for certain subspaces of $L^2[0, \infty)$, we prefer to concern ourselves with the rich class of relations satisfied when $\alpha = 0$.

Recall that the $Gau\beta$ -Heine symbol, also known as the q-factorial (Gasper & Rahman, 1990), reads

$$(z;q)_0 = 1,$$
 $(z;q)_n = (1-q^{n-1}z)(z;q)_{n-1} = \prod_{j=0}^{n-1} (1-q^j z),$ $j = 1, 2, \dots$

Since

$$\prod_{k=0}^{j-1} (q^j + q^k) = q^{(j-1)j/2} (-q;q)_j,$$

$$\prod_{k=j}^{m-1} (q^j + q^k) = q^{(m-j)j} (-1;q)_{m-j},$$

$$\begin{split} &\prod_{\ell=0}^{j-1} (q^j - q^\ell) \; = \; (-1)^j \, q^{(j-1)j/2}(q;q)_j, \qquad \text{and} \\ &\prod_{\ell=j+1}^m (q^j - q^\ell) \; = \; q^{(m-j)j}(q;q)_j, \end{split}$$

substitution into (2.2) results in the explicit form

$$R_m(t) := r_{m+1}^{[0]}(t) = \sum_{j=0}^m (-1)^j \frac{(-q;q)_j (-1;q)_{m-j}}{(q;q)_j (q;q)_{m-j}} e^{-q^j t}, \qquad m = 0, 1, \dots$$
(2.3)

 Let

$$G(t,z) := \sum_{m=0}^{\infty} R_m(t) z^m, \qquad |z| < 1.$$

Proposition 9 The function G obeys the functional-differential equation

$$\frac{\partial}{\partial t}G(t,z) + G(t,z) = -z \left[G(qt,z) - q\frac{\partial}{\partial t}G(qt,z)\right]$$
(2.4)

with the initial condition

$$G(0,z) = \frac{1}{1-z}.$$
(2.5)

 $\mathit{Proof}~$ We multiply (2.3) by z^m and sum. Interchanging the order of summation, we obtain

$$G(t,z) = \sum_{m=0}^{\infty} \sum_{j=0}^{m} (-1)^{j} \frac{(-q;q)_{j}(-1;q)_{m-j}}{(q;q)_{j}(q;q)_{m-j}} e^{-q^{j}t} z^{m}$$

$$= \sum_{j=0}^{\infty} (-1)^{j} \frac{(-q;q)_{j}}{(q;q)_{j}} e^{-q^{j}t} \sum_{m=j}^{\infty} \frac{(-1;q)_{m-j}}{(q;q)_{m-j}} z^{m}$$

$$= \left[\sum_{j=0}^{\infty} (-1)^{j} \frac{(-q;q)_{j}}{(q;q)_{j}} e^{-q^{j}t} z^{j} \right] \times \left[\sum_{m=0}^{\infty} \frac{(-1;q)_{m}}{(q;q)_{m}} z^{m} \right]$$

However, according to the Heine formula for basic hypergeometric functions (Gasper & Rahman, 1990),

$$\sum_{m=0}^{\infty} \frac{(-1;q)_m}{(q;q)_m} z^m = {}_1 \Phi_0 \left[\begin{array}{c} -1; \\ --; \end{array} q; z \right] = \frac{(-z;q)_{\infty}}{(z;q)_{\infty}},$$

therefore

$$G(t,z) = \frac{(-z;q)_{\infty}}{(z;q)_{\infty}} \sum_{j=0}^{\infty} (-1)^j \frac{(-q;q)_j}{(q;q)_j} e^{-q^j t} z^j,$$

and this, according to (Iserles, 1993), is the Dirichlet series expansion of the solution of the *pantograph equation* (2.4).

To evaluate the initial condition we again sum a basic hypergeometric series with the Heine formula,

$$\sum_{j=0}^{\infty} (-1)^j \frac{(-q;q)_j}{(q;q)_j} z^j = {}_1 \Phi_0 \left[\begin{array}{c} -q; \\ --; \end{array} q; z \right] = \frac{(qz;q)_{\infty}}{(-z;q)_{\infty}}.$$

Therefore

$$G(0,z) = \frac{(-z;q)_{\infty}}{(z;q)_{\infty}} \times \frac{(qz;q)_{\infty}}{(-z;q)_{\infty}} = \frac{1}{1-z},$$

affirming (2.5).

Corollary 10 The Taylor expansion (in t) of the function G is

$$G(t,z) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \frac{(-z;q)_k}{(z;q)_{k+1}} t^k.$$
(2.6)

Proof According to (Iserles, 1993), the solution of the pantograph equation

 $y'(t) = ay(t) + by(qt) + cy'(qt), \quad t \ge 0, \qquad y(0) = y_0, \tag{2.7}$

where $a, b, c \in \mathbb{C}$, $a \neq 0$, |c| < 1, can be expanded into the Taylor series

$$y(t) = y_0 \sum_{k=0}^{\infty} \frac{1}{k!} \frac{(-b/a;q)_k}{(c;q)_k} (at)^k.$$

Letting a = -1, b = -z, c = qz and $y_0 = 1/(1-z)$ yields (2.6).

We note as an aside that, according to (Iserles, 1993) and (Iserles & Liu, 1994), the solution of (2.7) exists and it is unique subject to the inequality |c| < 1. Moreover, $\operatorname{Re} a < 0$, |a| > |b| implies that the solution is asymptotically stable. Thus, thanks to the restriction |z| < 1, we deduce that $\lim_{t\to\infty} G(t, z) = 0$.

Next we consider recurrence relations that are obeyed by the sequence $\{R_n\}_{n=0}^{\infty}$. A differential recurrence follows by substitution in (1.11), namely

$$R'_{n-1} - R'_n = q^{n-1}R_{n-1} + q^n R_n, \qquad n = 1, 2, \dots$$
(2.8)

This recurrence can be rewritten in the form (1.12).

Several other relations can be derived by a moderate effort. For example, we can express R_n in terms of R_{n-1} and its derivatives at t and qt (a differential recurrence with rescaling). Thus, let

$$\varphi_n(t) := (1-q^n) R_n(t) + R_{n-1}(qt) - q^{n-1} R_{n-1}(t) + R'_{n-1}(t) - q R'_{n-1}(qt), \quad n = 1, 2, \dots$$
(2.9)

It is a trite exercise to verify that

$$\langle e^{-q^{\ell}t}, R_{n-1}(qt) \rangle = q^{-1} \langle e^{-q^{\ell-1}t}, R_{n-1} \rangle, \langle e^{-q^{\ell}t}, R'_{n-1}(qt) \rangle = -q^{-1} + q^{\ell-2} \langle e^{-q^{\ell-1}t}, R_{n-1} \rangle,$$
 $\ell = 1, 2, \dots,$

consequently

$$\langle e^{-q^{\ell}t}, \varphi_{n} \rangle = (1-q^{n}) \langle e^{-q^{\ell}t}, R_{n} \rangle + q^{-1} \langle e^{-q^{\ell-1}t}, R_{n-1} \rangle - q^{n-1} \langle e^{-q^{\ell}t}, R_{n-1} \rangle + q^{\ell} \langle e^{-q^{\ell}t}, R_{n-1} \rangle - q^{\ell-1} \langle e^{-q^{\ell-1}t}, R_{n-1} \rangle, \qquad \ell = 1, 2, \dots$$

Recalling that each R_m is orthogonal to $\exp(-q^{\ell}t)$ for $\ell = 0, 1, \ldots, m-1$, we thus deduce the relations

$$\langle \mathrm{e}^{-q^{\ell}t}, \varphi_n \rangle = 0, \qquad \ell = 0, 1, \dots, n-1.$$

Therefore φ_n is a scalar multiple of R_n . But Corollary 4 implies $R_n(0) = 1$, whilst (2.8) yields $R'_n(0) = -(2 - q^n - q^{n+1})/(1 - q)$. Substitution in (2.9) verifies at once $\varphi_n(0) = 0$, thus leading to the conclusion that $\varphi_n \equiv 0$.

Proposition 11 The sequence $\{R_n\}_{n=0,1,...}$ satisfies the differential recurrence relation with rescaling

$$(1-q^n)R_n(t) = q^{n-1}R_{n-1}(t) - R_{n-1}(qt) - R'_{n-1}(t) + qR'_{n-1}(qt), \qquad n = 1, 2, \dots$$

Finally, we report a pure recurrence relation with rescaling.

Proposition 12 For every $n = 2, 3, \ldots$ we have the relation

$$(1-q^n)R_n(t) = (1+q^{n-1})R_{n-1}(t) - (1+q^n)R_{n-1}(qt) + (1-q^{n-1})R_{n-2}(qt).$$
(2.10)

Proof Although (2.10) can be proved by comparing coefficients in (2.3), it is perhaps more interesting to use the generating function G. Comparing the coefficients in (2.6) affirms the identity

$$(1-z)[G(t,z) + zG(qt,z)] = (1+z)[G(t,gz) - qzG(qt,qz)],$$

which we rewrite in the form

$$G(t, z) - G(t, gz) = z[G(t, z) + G(t, qz)] - z[G(qt, z) - qG(qt, gz)] + z^{2}[G(qt, z) - qG(qt, qz)].$$

The recurrence (2.10) follows at once, by substituting the definition of the generating function G.

3. Convergence of projections

Let $\lambda_0, \lambda_1, \ldots$ be given positive numbers. We consider the approximation of a function $f \in L^2[0,\infty)$ by projections onto each H_n . We henceforth consider only the inner product $\langle f, g \rangle = \int_0^\infty f(t)g(t) \, dt$, although our discussion can be generalized with minimal effort to the inner product (2.1).

The orthogonal projection of f onto H_n can be written as a generalized Fourier series, which in turn can be expressed explicitly in terms of the Fourier transform of f, so that

$$F_n(t) = \sum_{m=0}^n \frac{\langle f, r_m \rangle}{\langle r_m, r_m \rangle} r_m(t) = 2 \sum_{m=0}^{n-1} \lambda_m \langle f, r_m \rangle r_m(t), \qquad n = 0, 1, \dots,$$
(3.1)

where (cf. (1.6))

$$r_m(t) = \sum_{\ell=0}^n r_{m,\ell} \exp(-\lambda_\ell t),$$

nplies $\langle f, r_m \rangle = \sum_{\ell=0}^m r_{m,\ell} \int_0^\infty f(t) e^{-\lambda_\ell t} dt = \sum_{\ell=0}^m r_{m,\ell} \hat{f}(-i\lambda_\ell).,$ (3.2)

in

The behaviour of the Fourier transform is entwined with the convergence of the projections. Let us use \mathcal{H} to denote the closure of the subspace $H_1 + H_2 + \cdots$ in the L^2 norm.

Proposition 13 Let $f \in \mathcal{H}$. If the function

$$\Lambda(z) := \sum_{n=0}^{\infty} \lambda_n z^n$$

is analytic in an open disc centred at the origin and we have the inequality

$$|\langle f, r_n \rangle| \le \omega^n |\langle f, r_0 \rangle|, \qquad n = 0, 1, \dots,$$
(3.3)

where $\omega \in (0, 1)$ is a point in the disc, then

$$||f - F_n||^2 \le 2|\langle f, r_0 \rangle|^2 \sum_{m=n}^{\infty} \lambda_m \omega^{2m},$$
 (3.4)

and the right hand side converges to zero as n tends to infinity.

Proof It follows from (3.1) by the Parseval equality that

$$||f - F_n||^2 = 2\sum_{m=n}^{\infty} \lambda_m |\langle f, r_m \rangle|^2,$$

and we deduce (3.4) from (3.3) and the analyticity of Λ at ω^2 .

Let us assume further that f is analytic in the closed unit disc. Substituting

$$\hat{f}(iz) = \sum_{k=0}^{\infty} \frac{\varphi_k}{k!} z^k$$

into (3.2) yields

$$\langle f, r_m \rangle = \sum_{k=0}^{\infty} \frac{\varphi_k}{k!} \sum_{\ell=0}^m r_{m,\ell} (-\lambda_\ell)^k = \sum_{k=0}^{\infty} \frac{\varphi_k}{k!} r_m^{(k)}(0), \qquad m = 0, 1, \dots$$
(3.5)

This formula and inequality (3.3) motivate our interest in the magnitude of the derivatives of r_m at the origin.

We restrict our attention in the remainder of this section to the parameters $\lambda_k = q^k$, $k = 0, 1, \ldots$, that have already featured in Section 2. According to (2.6) we have

$$\sum_{n=0}^{\infty} R_n^{(k)}(0) z^n = \frac{\partial^k}{\partial t^k} G(0, z) = (-1)^k \frac{(-z; q)_k}{(z; q)_{k+1}}, \qquad k = 0, 1, \dots,$$

hence the recursion

$$\begin{split} \sum_{n=0}^{\infty} R_n^{(k)}(0) z^n &= -\frac{1+q^{k-1}z}{1-q^k z} \sum_{n=0}^{\infty} R_n^{(k-1)}(0) z^n \\ &= -\sum_{n=0}^{\infty} \left[\sum_{\ell=0}^n R_{n-\ell}^{(k-1)}(0) q^{k\ell} \right] z^n - \sum_{n=1}^{\infty} \left[\sum_{\ell=1}^n R_{n-\ell}^{(k-1)}(0) q^{k\ell-1} \right] z^n. \end{split}$$

We deduce that

$$R_n^{(k)}(0) = -\sum_{\ell=0}^n R_{n-\ell}^{(k-1)}(0) q^{k\ell} - q^{k-1} \sum_{\ell=0}^{n-1} R_{n-1-\ell}^{(k-1)}(0) q^{k\ell}, \qquad n = 0, 1, \dots, \ k = 1, 2, \dots.$$
(3.6)

Proposition 14 The derivatives of R_n at the origin obey the inequality

$$|R_n^{(k)}(0)| \le \frac{(-1;q)_k}{(q;q)_k}, \qquad n,k = 0,1,\dots.$$
(3.7)

Proof We use induction on the derivative order k. The inequality is true for k = 0 because, by Corollary 4, $R_n(0) = 1$. We thus assume its correctness for k - 1 and employ (3.6) to argue that

$$\begin{aligned} |R_n^{(k)}(0)| &\leq \frac{(-1;q)_{k-1}}{(q;q)_{k-1}} \left[\sum_{\ell=0}^n q^{k\ell} + q^{k-1} \sum_{\ell=0}^{n-1} q^{k\ell} \right] \\ &\leq \frac{(-1;q)_{k-1}}{(q;q)_{k-1}} \left[\frac{1}{1-q^k} + \frac{q^{k-1}}{1-q^k} \right] = \frac{(-1;q)_k}{(q;q)_k} \end{aligned}$$

which completes the proof.

Corollary 15

$$|R_n^{(k)}(0)| \le \frac{(-1;q)_{\infty}}{(q;q)_{\infty}}, \qquad n,k = 0,1,\dots.$$
(3.8)

Proof This is immediate from (3.7), since the sequence

$$\left\{\frac{(-1;q)_{k-1}}{(q;q)_{k-1}}: k = 0, 1, \ldots\right\}$$

is increasing for every $q \in (0, 1)$.

Of course, analyticity of \hat{f} in the closed unit disc implies absolute convergence of its Taylor series at z = 1. Setting

$$\sigma := \sum_{k=0}^{\infty} \frac{|\varphi_k|}{k!} < \infty,$$

we thereby deduce from (3.8) that

$$|\langle f, R_n \rangle \le \sigma \frac{(-1;q)_{\infty}}{(q;q)_{\infty}}, \qquad n = 0, 1, \dots$$

Consequently, (3.5) and (3.8) imply the bound

$$|\langle f, R_n \rangle| \le \sigma \frac{(-1;q)_{\infty}}{(q;q)_{\infty}}, \qquad n = 0, 1, \dots,$$

and, when $f \in \mathcal{H}$, the Parseval theorem yields the expression

$$||f - F_N||^2 = 2 \sum_{n=N}^{\infty} q^n |\langle f, R_n \rangle|^2$$

Therefore, our bound on $|\langle f, R_n \rangle|$ establishes the following result.

Theorem 16 If $f \in \mathcal{H}$ and \hat{f} is analytic in the closed unit disc, then

$$\|f - F_n\| \le \sigma^* q^{N/2} \xrightarrow{N \to \infty} 0, \tag{3.9}$$

where

$$\sigma^* = \sigma \sqrt{\frac{2}{1-q}} \times \frac{(-1;q)_{\infty}}{(q;q)_{\infty}}.$$

Furthermore, since

$$\sum_{k=0}^{\infty} \frac{|\varphi_k|}{k!} \le \int_0^{\infty} e^t |f(t)| \, \mathrm{d}t.$$
(3.10)

Therefore boundedness of this integral implies analyticity of f in the closed unit disc. Let us return to the interesting special case $f(t) = \exp(-\lambda t)$. By (1.2),

$$\langle e^{-\lambda t}, R_n \rangle = \langle g_\lambda, R_n \rangle = R_n(-i\lambda)$$

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and (1.5) yields

$$\langle e^{-\lambda t}, r_n \rangle = \frac{\prod_{k=0}^{n-1} (\lambda - \lambda_k)}{\prod_{k=0}^{n} (\lambda + \lambda_k)}, \qquad n = 0, 1, \dots$$

Specialising to $\lambda_{\ell} = q^{\ell}$, we obtain

$$\langle e^{-\lambda t}, R_n \rangle = \frac{1}{\lambda} \frac{(1/\lambda; q)_n}{(-1/\lambda; q)_{n+1}}, \qquad n = 0, 1, \dots$$
 (3.11)

We thus deduce by the method of proof of Theorem 16 that

$$F_n(t) = \frac{2}{\lambda} \sum_{n=0}^{N} \frac{(1/\lambda; q)_n}{(-1/\lambda; q)_{n+1}} q^n R_n(t)$$
(3.12)

converges in norm to the orthogonal projection of $\exp(-\lambda t)$ on \mathcal{H} .

The projection of $f(t) = \exp(-\lambda t)$, $\lambda > 1$, onto H_n is

$$F_N(t) = \frac{2}{\lambda} \sum_{n=0}^{N} \frac{(1/\lambda; q)_n}{(-1/\lambda; q)_{n+1}} q^n R_n(t).$$

Letting $N \to \infty$, it is easy to verify that

$$F_N(t) \to F(t) = \frac{2}{1+\lambda} \sum_{n=0}^{\infty} \frac{(1/\lambda; q)_n}{(-q/\lambda; q)_n} q^n R_n(t).$$
 (3.13)

We next substitute the explicit expression for R_n from (2.3), whence exchanging the order of summation yields

$$F(t) = \frac{2}{1+\lambda} \sum_{m=0}^{\infty} (-1)^m \frac{(-q;q)_m}{(q;q)_m} e^{-q^m t} \sum_{n=m}^{\infty} \frac{(1/\lambda;q)_n (-1;q)_{n-m}}{(-q/\lambda;q)_n (q;q)_{n-m}} q^n$$

$$= \frac{2}{1+\lambda} \sum_{m=0}^{\infty} (-1)^m \frac{(-q;q)_m}{(q;q)_m} q^m e^{-q^m t} \sum_{n=0}^{\infty} \frac{(1/\lambda;q)_{m+n} (-1;q)_n}{(q;q)_n (-q/\lambda;q)_{m+n}} q^n$$

$$= \frac{2}{1+\lambda} \sum_{m=0}^{\infty} (-1)^m \frac{(-q;q)_m (1/\lambda;q)_m}{(q;q)_m (-q/\lambda;q)_m} q^m e^{-q^m t} \sum_{n=0}^{\infty} \frac{(q^m/\lambda;q)_n (-1;q)_n}{(q;q)_n (-q^{m+1}/\lambda;q)_n} q^n$$

$$= \frac{2}{1+\lambda} \sum_{m=0}^{\infty} (-1)^m \frac{(-q;q)_m (1/\lambda;q)_m}{(q;q)_m (-q/\lambda;q)_m} q^m e^{-q^m t} 2\Phi_1 \begin{bmatrix} q^m/\lambda, -1; \\ -q^{m+1}/\lambda; q,q \end{bmatrix} (3.14)$$

– we refer to (Gasper & Rahman, 1990) for the terminology of basic hypergeometric series.

The $_2\Phi_1$ series in (3.14) can be summed with the q-Gauß formula (Gasper & Rahman, 1990, formula II.8, page 236),

$${}_{2}\Phi_{1}\left[\begin{array}{c}q^{m}/\lambda,-1;\\-q^{m+1}/\lambda;\end{array} q,q\right] = \frac{(-q;q)_{\infty}(q^{m+1}/\lambda;q)_{\infty}}{(q;q)_{\infty}(-q^{m+1}/\lambda;q)_{\infty}}.$$

Substitution in (3.14) and elementary simplification result in

$$F(t) = \frac{1}{\lambda} \frac{(-1;q)_{\infty} (1/\lambda;q)_{\infty}}{(q;q)_{\infty} (-1/\lambda;q)_{\infty}} \sum_{m=0}^{\infty} (-1)^m \frac{(-q;q)_m}{(q;q)_m} \frac{q^m}{1 - q^m/\lambda} e^{-q^m t}.$$
 (3.15)

Note that (3.15) is valid for all $\lambda > 0$, $\lambda \neq q^k$ for $k \ge 0$. In the case $\lambda = q^k$ it is easy to prove that $F(t) = \exp(-q^k t)$.

Finally, we explicitly evaluate F(0) from (3.15). In the course of our analysis we twice use an explicit formula for the summation of $_{1}\Phi_{0}$ series (Gasper & Rahman, 1990, formula II.3, page 236).

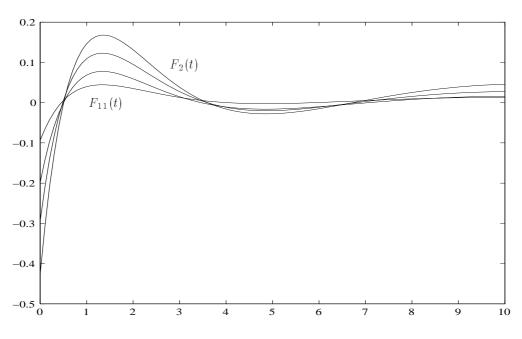


Figure 1 The functions F_N for $N = 2, 5, 8, 11, q = \frac{1}{2}$ and $\lambda = \frac{3}{2}$.

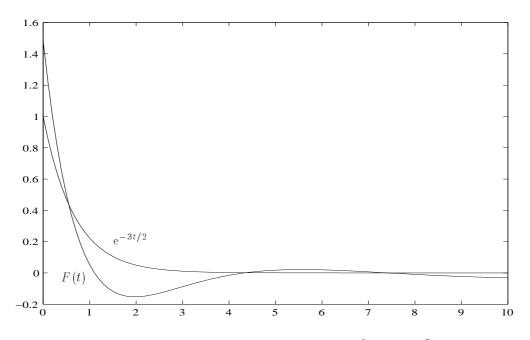


Figure 2 The functions F and $\exp(-\lambda t)$ for $q = \frac{1}{2}$ and $\lambda = \frac{3}{2}$.

Bearing in mind that $0 < 1/\lambda < 1$, we expand into series,

$$\begin{split} &\sum_{m=0}^{\infty} (-1)^m \frac{(-q;q)_m}{(q;q)_m} \frac{q^m}{1-q^m/\lambda} \\ &= \sum_{m=0}^{\infty} (-1)^m \frac{(-q;q)_m}{(q;q)_m} q^m \sum_{\ell=0}^{\infty} \frac{q^{m\ell}}{\lambda^{\ell}} \\ &= \sum_{\ell=0}^{\infty} \frac{1}{\lambda^{\ell}} {}_1 \Phi_0 \left[-q; -; q, -q^{\ell+1} \right] \\ &= \sum_{\ell=0}^{\infty} \frac{1}{\lambda^{\ell}} \frac{(q^{\ell+2};q)_{\infty}}{(-q^{\ell+1};q)_{\infty}} = \lambda \left[\sum_{\ell=0}^{\infty} \frac{(q^{\ell+1};q)_{\infty}}{(-q^{\ell};q)_{\infty}} \frac{1}{\lambda^{\ell}} - \frac{(q;q)_{\infty}}{(-1;q)_{\infty}} \right] \\ &= \lambda \frac{(q;q)_{\infty}}{(-1;q)_{\infty}} \left[\sum_{\ell=0}^{\infty} \frac{(-1;q)_{\ell}}{(q;q)_{\ell}} \lambda^{-\ell} - 1 \right] \\ &= \lambda \frac{(q;q)_{\infty}}{(-1;q)_{\infty}} \left\{ {}_1 \Phi_0 \left[\begin{array}{c} -1; \\ -; \end{array}; q, \lambda^{-1} \right] - 1 \right\} = \frac{(q;q)_{\infty}}{(-1;q)_{\infty}} \left[\frac{(-1/\lambda;q)_{\infty}}{(1/\lambda;q)_{\infty}} - 1 \right]. \end{split}$$

Therefore, substitution in (3.15) proves that

$$F(0) = 1 - \frac{(1/\lambda; q)_{\infty}}{(-1/\lambda; q)_{\infty}}.$$

Note that, unless $\lambda = q^{-m}$ for a nonnegative integer m, it follows that $F(0) \neq 1$. This constitutes a formal proof of the statement that the sequence $\{F_N\}$ does not converge to $\exp(-\lambda t)$; of course we have already provided a stronger result in Proposition 7. We can also use the infinite product 1.9 to characterize the limit of the projections $\{F_N\}$. Indeed, the proof of Proposition 7 implies that

$$\lim_{N \to \infty} \hat{f}(z) - \hat{F}_N(z) = \frac{-i}{z - i\lambda} \Delta(iz) \Delta(\lambda).$$
(3.16)

In Figure 1 we display the functions F_N for different values of N in the case $q = \frac{1}{2}$, $\lambda = \frac{3}{2}$. It illustrates vividly our observation that $F_N \to F$ yet, as can be seen in Figure 2, the function F is distinct from $\exp(-\frac{3}{2}t)$.

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